# 945. Molecular Polarisability. Specificution of the Phosphite, Phosphate, Thiophosphate, and Arsenite Group Polarisability Ellipsoid Semi-axes. 

By M. J. Aroney, R. J. W. Le Fèvre, and J. Saxby.

On the basis of polarisability ellipsoids of revolution, major and minor semi-axes for groups $\mathrm{O}_{3} \mathrm{X}$, where $\mathrm{X}=\mathrm{P}, \mathrm{PO}, \mathrm{PS}$, or As, appear as 0.410 and $0.359,0.416$ and $0.284,0.869$ and 0.492 , and 0.535 and 0.421 (all in $10^{-23}$ c.c. units), respectively. "Effective" longitudinal and transverse polarisabilities of the $\mathrm{O}-\mathrm{P}$ and $\mathrm{O}-$ As bonds are estimated as $b_{\mathrm{L}}^{0-\mathrm{P}}=0.086$, $b_{\mathrm{T}}^{\mathrm{O}-\mathrm{P}}=0.145, b_{\mathrm{L}}^{\mathrm{O}-\mathrm{As}}=0.066$, and $b_{\mathrm{T}}^{\mathrm{O}-\mathrm{As}}=0.196$ (also in $10^{-23}$ c.c. units).
OUR investigation of the configurations of molecules (RO) ${ }_{3} \mathrm{X}$ (where X is $\mathrm{P}, \mathrm{PO}, \mathrm{PS}$, and As), as solutes in non-polar media, requires a knowledge of the anisotropic polarisabilities of each bond or group in the system. Accordingly the present work has been undertaken to determine the polarisability specifications of the groups $\mathrm{O}_{3} \mathrm{P}, \mathrm{O}_{3} \mathrm{PO}, \mathrm{O}_{3} \mathrm{PS}$, and $\mathrm{O}_{3} \mathrm{As}$, to supplement previous data ${ }^{1,2}$ for the bonds $\mathrm{C}-\mathrm{C}, \mathrm{C}-\mathrm{O}$, and $\mathrm{C}-\mathrm{H}\left(b_{\mathrm{L}}^{\mathrm{C}-\mathrm{C}}=0.099, b_{\mathrm{T}}^{\mathrm{C}-\mathrm{C}}=\right.$ $\left.b_{\mathrm{V}}^{\mathrm{C}-\mathrm{C}}=0.027 ; b_{\mathrm{L}}^{\mathrm{C}-\mathrm{O}}=0.090, b_{\mathrm{T}}^{\mathrm{C}-\mathrm{O}}=b_{\mathrm{V}}^{\mathrm{C}-\mathrm{O}}=0.043 ; b_{\mathrm{L}}^{\mathrm{C}-\mathrm{H}}=b_{\mathrm{T}}^{\mathrm{C}-\mathrm{H}}=b_{\mathrm{V}}^{\mathrm{C}-\mathrm{H}}=0.064\right) . *$

## Experimental

Materials, Apparatus, ctc.-The solutes were prepared and purified by the method of Verkade and Reynolds. ${ }^{3}$ Apparatus, techniques, symbols used, and methods of calculation have been described before. ${ }^{1,4}$ Measurements, taken in all cases at $25^{\circ}$,
 are listed in Table 1; $\Delta B$ is the difference found between the Kerr constant of the solvent and those of solutions containing weight fractions $w_{2}$ of solute. Estimates of $\delta\left(=\sum 10^{7} \Delta B / B_{1} \sum w_{2}\right)$ deduced from Table 1 are given in Table 2 together with the various other data required for the calculation of the molar Kerr constants listed therein. The following constants apply at $25^{\circ}$ to dioxan: $\varepsilon_{1}=2.2090 ; d_{1}=1.0280 ;\left(n_{1}\right)_{\mathrm{D}}=$ $1.4202 ; 10^{2} B_{1}=0.068 ; 10^{12}{ }_{\mathrm{s}} K_{1}=0.0116$.

* Polarisabilities are expressed throughout this paper as $10^{-23}$ c.c. units.
${ }^{1}$ Le Fèvre and Le Fèvre, Rev. Pure Appl. Chem., 1955, 5, 261.
${ }^{2}$ Le Fèvre, unpublished data.
${ }^{3}$ Verkade and Reynolds, J. Org. Chem., 1960, 25, 663.
${ }^{4}$ Le Fèvre and Le Fèvre, Chap. XXXVI in " Physical Methods of Organic Chemistry," ed. Weissberger, Interscience Publ., Inc., New York, London, 3rd edn., Vol. I, p. 2459.

Table 1.
Incremental Kerr constants of solutions in dioxan at $25^{\circ}$.

| 4-Methyl-1-phospha-2,6,7-trioxabicyclo[2,2,2]octane ( $\mathrm{I} ; \mathrm{X}=\mathrm{P}$ ). |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $10^{6} w_{2}$ |  | 2831 | 3060 | 3420 | 3528 | 3997 | 44.72 | 4505 |
| $10^{7} \Delta B$ |  | $0 \cdot 043$ | 0.048 | 0.051 | $0 \cdot 054$ | $0 \cdot 059$ | $0 \cdot 066$ | 0.065 |
| whence $\sum 10^{7} \Delta B / \sum w_{2}=15 \cdot 0$. |  |  |  |  |  |  |  |  |
| 4-Methyl-1-phospha-2,6,7-trioxabicyclo [2,2,2]octane 1-oxide ( $\mathrm{I} ; \mathrm{X}=\mathrm{PO}$ ). |  |  |  |  |  |  |  |  |
| $10^{6} w_{2}$ |  |  | 672 | 704 | 762 | 763 | 881 | 967 |
| $10^{7} \Delta B$ |  |  | 0.037 | 0.037 | 0.042 | 0.043 | 0.050 | 0.054 |
| whence $\Sigma 10^{7} \Delta B / \Sigma w_{2}=55 \cdot 4$. |  |  |  |  |  |  |  |  |
| 4-Methyl-1-phospha-2,6,7-trioxabicyclo[2,2,2]octane 1-sulphide ( $\mathrm{I} ; \mathrm{X}=\mathrm{PS}$ ) |  |  |  |  |  |  |  |  |
| $10^{6} w_{2}$ | ...... | 562 | 644 | 645 | 682 | 714 | 761 | 829 |
| $10^{7} \Delta B$ | ...... | 0.048 | 0.054 | 0.055 | 0.057 | 0.064 | $0 \cdot 068$ | 0.073 |
| whence $\Sigma 10^{7} \Delta B / \Sigma w_{2}=86 \cdot 6$. |  |  |  |  |  |  |  |  |
| 4-Methyl-1-arsa-2,6,7-trioxabicyclo [2,2;2]octane ( $\mathrm{I} ; \mathrm{X}=\mathrm{As}$ ) |  |  |  |  |  |  |  |  |
| $10^{5} w_{2}$ |  |  | 1049 | 1594 | 1740 | 2189 | 2209 | 2501 |
| $10^{7} \Delta B$ |  |  | 0.050 | 0.090 | 0.093 | $0 \cdot 113$ | 0.096 | $0 \cdot 116$ |
| whence $\Sigma 10^{7} \Delta B / \Sigma w_{2}=4.95 . *$ |  |  |  |  |  |  |  |  |

* Solutions of this substance, when placed in the Kerr cell, quickly became yellow, so that the measurements of $\Delta B$ were made very rapidly; this results in a greater uncertainty in $\delta B_{1}$ than for the other solutes.

Table 2.
Polarisations, dipole moments, and molar Kerr constants.

| Solute | $\alpha \varepsilon_{1}{ }^{*}$ | $\beta \dagger$ | $\gamma \ddagger$ | $\delta$ | ${ }_{\mathbf{E}} P$ (c.c.) § | $\mu(\mathrm{D}){ }^{\text {* }}$ | $10^{12} \infty\left({ }_{m} K_{2}\right)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{CH}_{3} \cdot \mathrm{C}<\left[\mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3}>\mathrm{P}$ | $14 \cdot 6$ | $0 \cdot 184$ | 0.034 | 220 | 33.6 | $4 \cdot 15$ | 368 |
| $\mathrm{CH}_{3} \cdot \mathrm{C} \leqslant\left[\mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3} \cdots \mathrm{PO}$ | $38 \cdot 6$ | 0.320 | 0.044 | 814 | $32 \cdot 4$ | $7 \cdot 10$ | 1519 |
| $\mathrm{CH}_{3} \cdot \mathrm{C} \leq\left[\mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3}>\mathrm{PS}$ | $32 \cdot 0$ | $0 \cdot 244$ | 0.054 | 1274 | 39.7 | $6 \cdot 77$ | 2637 |
| $\mathrm{CH}_{3} \cdot \mathrm{C} \leq\left[\mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3}>\mathrm{As}$ | $3 \cdot 72$ | $0 \cdot 350$ | 0.042 | 73 | $35 \cdot 7$ | $2 \cdot 36$ | 160 |

* Quoted from Brown, Verkade, and Piper (J. Phys. Chem., 1961, 65, 2051). † Calc. from data of Brown et al. (loc. cit.). $\ddagger$ Calc. from $R_{\mathrm{D}}$ (cf. Brown et al., loc. cit.) and $\beta$. § Extrapolation of the refractivities listed by Vogel and Cowan ( $J ., 1943,16$ ) for trimethyl phosphate leads to $\infty R\left(\mathrm{Me}_{3} \mathrm{PO}_{4}\right)$ $={ }_{\mathrm{E}} P\left(\mathrm{Me}_{3} \mathrm{PO}_{4}\right)=27.4$ c.c., from which we obtain, after addition of four ${ }_{\mathrm{E}} P(\mathrm{C}-\mathrm{C})$ equivalents (cf. Le Fèvre and Steel, Chem. and Ind., 1961, 670), ${ }_{\mathbf{E}} P\left(\mathrm{CH}_{3} \cdot \mathrm{C}<\left[\mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3}>\mathrm{PO}\right)=32 \cdot \mathbf{4}$ c.c. Gillis [Rev. Pure Appl. Chem. (Australia), 1960,10,21] gives $R_{\mathrm{D}}(\mathrm{P}=\mathrm{O})=-1.22$ c.c. and $R_{\mathrm{D}}(\mathrm{P}=\mathrm{S})=6.4$ c.c., from which we estimate the electronic polarisation of the molecules $\mathrm{CH}_{3} \cdot \mathrm{C}<\left[\mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3}>\mathrm{P}$ and $\mathrm{CH}_{3} \cdot \mathrm{C}^{\prime}<\left[\mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3}>\mathrm{PS}$ as 33.6 and 39.7 c.c., respectively [assuming that ${ }_{\mathrm{E}} P(\mathrm{P}=\mathrm{O})=0.95 R_{\mathrm{D}}(\mathrm{P}=\mathrm{O})$ and that ${ }_{\mathrm{E}} P(\mathrm{P}=\mathrm{S})=$ $0.95 R_{\mathrm{D}}(\mathrm{P}=\mathrm{S})$ ]. From the dispersion data of Gryszkiewicz-Trochimowski and Sikorski (Bull. Soc. chim. France, 1927, 41, 1570), $\mathbf{E} P\left[(\mathrm{MeO})_{\mathbf{3}} \mathrm{As}\right]$ is calculated as $30 \cdot 7$ c.c.; addition to this of four ${ }_{\mathbf{E}} P(\mathrm{C}-\mathrm{C})$ equivalents yields ${ }_{\mathrm{E}} P\left(\mathrm{CH}_{3} \cdot \mathrm{C}\left[\mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3} \mathrm{As}\right)=35 \cdot 7$ c.c.


## Discussion

The ellipsoid of polarisability for each of the molecules $\left.\mathrm{CH}_{3} \cdot \mathrm{C}^{2} \mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3} \mathrm{X}$ can be regarded as one of revolution with $b_{1} \neq b_{2}=b_{3}$, i.e., $b_{1}$ is located along the symmetry axis (see I); it follows that $\mu_{1}=\mu_{\mathrm{obs}}, \mu_{2}=\mu_{3}=0$. The equations relating the electronic polarisation and molar Kerr constant with the principal polarisabilities simplify to:

$$
\begin{gather*}
{ }_{\mathrm{E}} P=4 \pi \boldsymbol{N}\left(b_{1}+2 b_{2}\right) / 9 ;  \tag{1}\\
{ }_{\mathrm{m}} K=4 \pi \boldsymbol{N}\left\{{ }_{\mathrm{D}} P\left(b_{1}-b_{2}\right)^{2} / \boldsymbol{k} T_{\mathrm{E}} P+\mu_{\mathrm{obs}}^{2}\left(b_{1}-b_{2}\right) / \boldsymbol{k}^{2} T^{2}\right\} / 405 . \tag{2}
\end{gather*}
$$

Substitution in equations (1) and (2) of $\mu_{\mathrm{obs}},{ }_{\mathrm{E}} P$, and ${ }_{\infty}\left({ }_{\mathrm{m}} K_{2}\right)$, from Table 2, and of ${ }_{\mathrm{D}} P$ (which we take as 1.05 times the corresponding $R_{\mathrm{D}}$ of ref. 5) leads to the molecular polarisability semi-axes listed in Table 3.

[^0]Table 3.
Polarisability semi-axes of the molecules $\mathrm{CH}_{3} \cdot \mathrm{C}<\left[\mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3}>\mathrm{X}$.

|  | $b_{1}$ | $b_{2}=b_{3}$ | $b_{1} / b_{2}$ |
| :---: | :---: | :---: | :---: |
| $\left.\mathrm{CH}_{3} \cdot \mathrm{C}^{2} \mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3} \mathrm{P}$ | 1.460 | 1.268 | $1 \cdot 15$ |
| $\mathrm{CH}_{3} \cdot \mathrm{C}^{\left(\mathrm{CH}_{3} \cdot \mathrm{O}\right]_{3} \mathrm{PO}}$ | $1 \cdot 466$ | $1 \cdot 193$ | $1 \cdot 23$ |
| $\mathrm{CH}_{3}{ }^{\circ} \mathrm{C}\left[\mathrm{CH}_{2}{ }^{\circ} \mathrm{O}\right]_{3} \mathrm{PS}$ | 1.919 | $1 \cdot 401$ | $1 \cdot 37$ |
| $\mathrm{CH}_{3}{ }^{-}\left[\mathrm{CH}_{2}{ }^{\circ} \mathrm{O}\right]_{3} \mathrm{As}$ | 1.585 | $1 \cdot 330$ | $1 \cdot 19$ |

Now $b_{1}$ and $b_{2}$ can be equated with the component bond and group polarisabilities through the following expressions:
and $\quad b_{1}+2 b_{2}=b_{1}^{0, \mathrm{x}}+2 b_{2}^{0_{2}^{\mathrm{x}}}+3 b_{\mathrm{L}}^{\mathrm{C}-0}+6 b_{\mathrm{T}}^{0-0}+4 b_{\mathrm{L}}^{\mathrm{C}} \mathrm{O}+8 b_{\mathrm{T}}^{0-0}+27 b_{\mathrm{L}}^{C-\mathrm{H}}$,
where $\phi$ and $\lambda$ are the angles of inclination of each $b_{\mathrm{L}}^{\mathrm{G}-0}$ and each non-terminal carboncarbon bond axis, respectively, with the $b_{1}$ direction. For the phosphorus compounds it seems reasonable to assume the following geometrical specifications: $r^{0-\mathrm{C}}=1.54 \AA$, $r^{-0}=1 \cdot 43 \AA, \angle C C C=\angle C C O=\angle C O P=$ tetrahedral, $r^{0-\mathrm{P}}=1 \cdot 57 \AA$, and $\angle \mathrm{OPO}=$ $106^{\circ}$ (from an $X$-ray study ${ }^{6}$ of phosphoric acid), from which we deduce $\phi=0^{\circ}$ and $\lambda=70^{\circ}$. Very few analogous data exist for the arsenite: $r^{0-A s}$ in $\mathrm{As}_{4} \mathrm{O}_{6}$ is given in ref. 7 as $1.78 \AA$, and it seems likely that $0<\phi<10^{\circ}$; the polarisability semi-axes given below for the $\mathrm{O}_{3} \mathrm{As}$ group have been calculated by using the same values of $\phi$ and $\lambda$ as for the phosphorus compounds.

Tables 2 and 3 show the following information: (a) The ratios $b_{1}: b_{2}$ are of the same order of magnitude, since the molecules are sterically very similar. (b) Electronic polarisation for the phosphate is smaller than for the phosphite, so that $\sum^{b}$ (phosphate) $<$ $\sum b$ (phosphite); the addition to the phosphite of an oxygen atom does not significantly change the polarisability along the $b_{1}$ direction: $\Delta b_{1}=b_{1}$ (phosphate) $-b_{1}$ (phosphite) $=$ $+0.006, \Delta b_{2}=-0.075$ (cf. Hacket and Le Fèvre's conclusions ${ }^{8}$ for phosphorus trichloride and oxychloride, where $\Delta b_{1}=-0.01 \dot{2}$ to +0.044 , and $\Delta b_{2}=-0.032$ to -0.060 ). (c) Addition of a sulphur atom to the phosphite results in an increase of polarisability directed predominantly along the symmetry axis: $\Delta b_{1}=+0 \cdot 459, \Delta b_{2}=+0.133$.

Verkade and Reynolds ${ }^{3}$ record that in the infrared spectra of these compounds the $\mathrm{P}=\mathrm{O}$ and $\mathrm{P}=\mathrm{S}$ stretching frequencies occur at 1325 and $800 \mathrm{~cm} .^{-1}$, respectively. Substitution of these values into the empirical equation ${ }^{9}$ of Le Fèvre (1959) together with the intercentre distances, $r^{\mathrm{P}=0}=1.45$ and $r^{\mathrm{P}=\mathrm{S}}=1.85$ (from Williams et al. ${ }^{10}$ ), leads to the following predicted longitudinal polarisabilities: $b_{\mathrm{L}}^{\mathrm{P}=\mathrm{O}}($ calc. $)=0 \cdot 484, b_{\mathrm{L}}^{\mathrm{P}=\mathrm{S}}($ calc. $)=0.928$. Significantly the difference ( $b_{\mathrm{L}}^{\mathrm{E}=\mathrm{s}}$ calc. $-b_{\mathrm{L}}^{\mathrm{P}=0}$ calc.) equals +0.444 in good agreement with that found from experiment $\left\{b_{1}\left(\mathrm{CH}_{3} \cdot \mathrm{C}<\left[\mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3}>\mathrm{PS}\right)-b_{1}\left(\mathrm{CH}_{3} \cdot \mathrm{C}<\left[\mathrm{CH}_{2} \cdot \mathrm{O}\right]_{3}>\mathrm{PO}\right)=\right.$ $+0 \cdot 453\}$.

Table 4.
Polarisability semi-axes of the groups $\mathrm{O}_{3} \mathrm{X}$.

|  | $\mathrm{O}_{3} \mathrm{P}$ | $\mathrm{O}_{3} \mathrm{PO}$ | $\mathrm{O}_{3} \mathrm{PS}$ | $\mathrm{O}_{3} \mathrm{As}$ |
| :---: | :---: | :---: | :---: | :---: |
| $b_{1} \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots$ | $0 \cdot 410$ | 0.416 | $0 \cdot 869$ | $0 \cdot 535$ |
| $b_{2}=b_{3} \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots$ | $0 \cdot 359$ | $0 \cdot 284$ | $0 \cdot 492$ | $0 \cdot 421$ |

The " effective" longitudinal and transverse polarisabilities of the $\mathrm{O}-\mathrm{P}$ and $\mathrm{O}^{-} \mathrm{As}$

[^1]bonds in the phosphite and arsenite can be deduced from the $\mathrm{O}_{3} \mathrm{P}$ and $\mathrm{O}_{3}$ As group values of Table 4 by means of the equations:
\[

$$
\begin{gather*}
b_{\mathrm{B}}^{0_{\mathrm{a}}^{\mathrm{a} \mathrm{X}}+2 b_{\mathrm{a}}^{0_{\mathrm{o}} \mathrm{X}}=3 b_{\mathrm{L}}^{0-\mathrm{x}}+6 b_{\mathrm{R}}^{0-\mathrm{x}},}  \tag{5}\\
b_{1}^{0, \mathrm{x}}=3 b_{\mathrm{L}}^{0-\mathrm{x}} \cos ^{2} \psi+3 b_{\mathrm{T}}^{0-\mathrm{x}} \sin ^{2} \psi,
\end{gather*}
$$
\]

and
where $\psi$, the angle which the longitudinal $\mathrm{O}^{-} \mathrm{P}$ or $\mathrm{O}-$ As axis makes with the $b_{1}$ direction, is taken as $68^{\circ}$ (by analogy with Furberg's data). Solution yields $b_{\mathrm{L}}^{0-\mathrm{P}}=0 \cdot 086, b_{\mathrm{T}}^{\mathrm{O}-\mathrm{P}}=$ $0 \cdot 145, b_{\mathrm{L}}^{\mathrm{O}-\mathrm{As}}=0.066, b_{\mathrm{T}}^{\mathrm{OAS}}=0 \cdot 196$. The "effective" bond semi-axes thus resolved include, of course, the polarisability contributions of the " lone-pair" electrons.


[^0]:    ${ }^{5}$ Brown, Verkade, and Piper, J. Phys. Chem., 1961, 65, 2051.
    7 U 2

[^1]:    ${ }^{6}$ Furberg, Acta Chem. Scand., 1955, 9, 1557.
    7 Sutton et al., " Tables of Interatomic Distances and Configuration in Molecules and Ions," Chem. Soc. Spec. Publ. No. 11, 1958.
    ${ }^{8}$ Hacket and Le Fèvre, $J ., 1961,2612$.
    ${ }^{9}$ Le Fèvre, Proc. Chem. Soc., 1959, 363.
    ${ }^{10}$ Williams, Sheridan, and Gordy, J. Chem. Phys., 1952, 20, 164.

